

All Electronic Phase-Retrieval Measurement of  
Femtosecond Optical Waveform

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## Outline

### Characterization of femtosecond optical waveform by nonlinear wave mixing

1. autocorrelator
2. frequency resolved optical gating

### All electronic phase retrieval measurement using nonlinear detectors

1. experimental setup
2. waveform reconstruction algorithm

### Potential applications

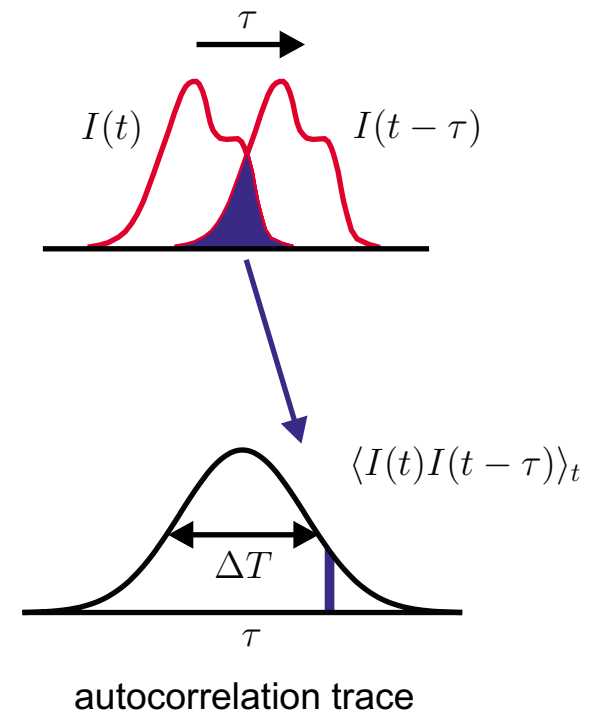
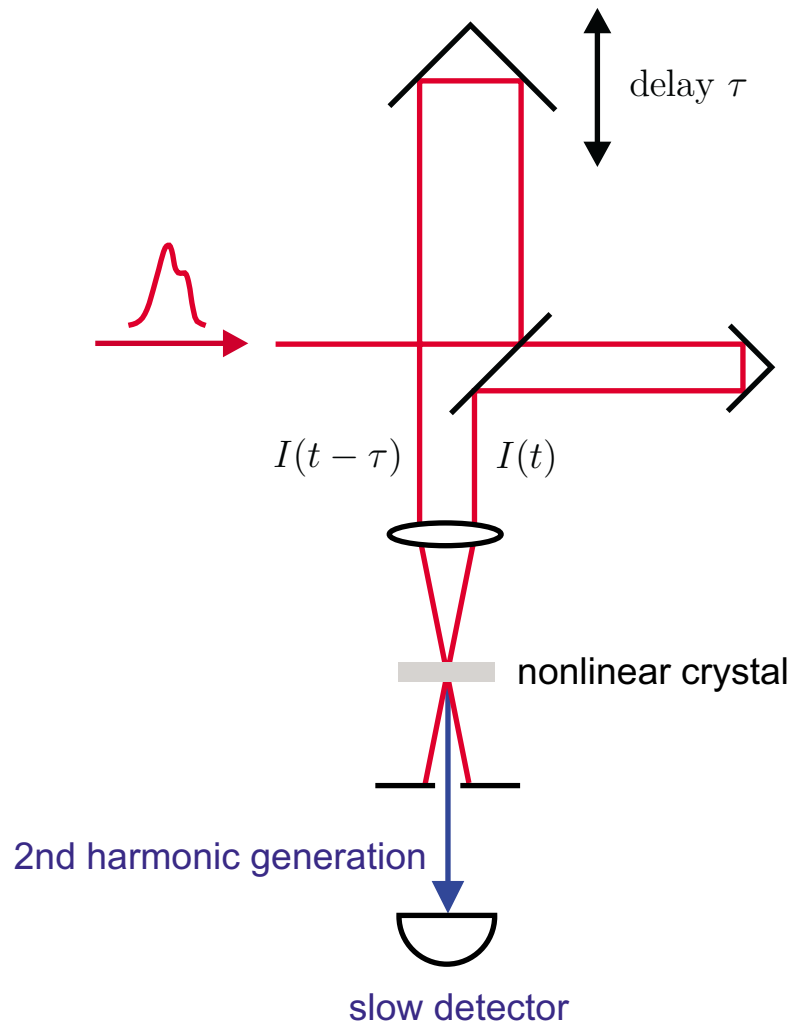
1. ultrafast spectroscopy
2. beam diagnosis for accelerator

# Characterization of femtosecond optical waveform by nonlinear wave mixing 📄

1. autocorrelator

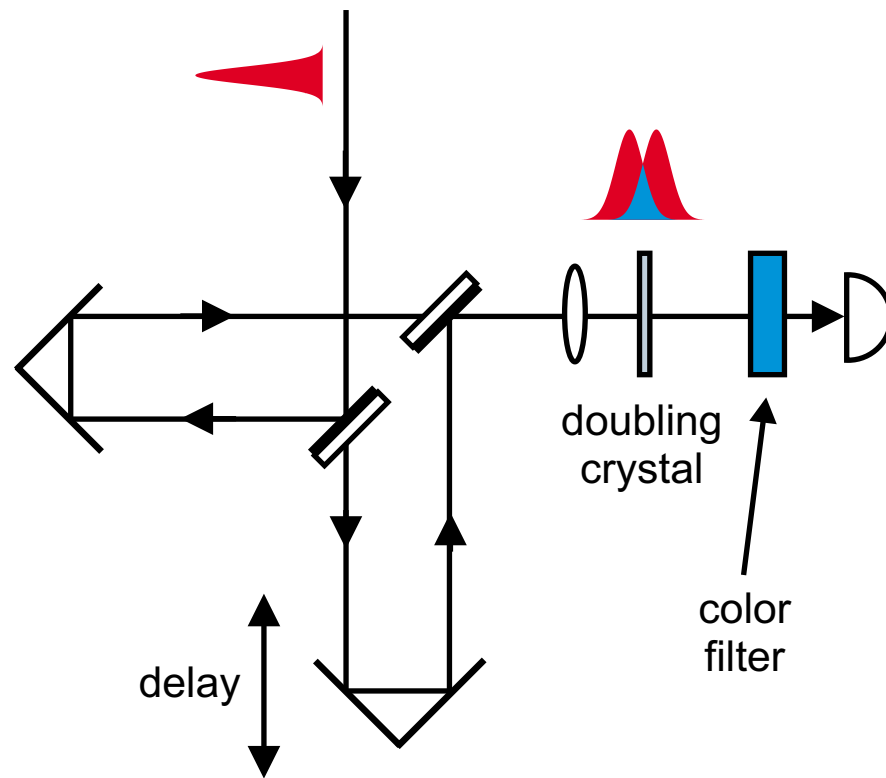
2. frequency resolved optical gating

# Intensity autocorrelator

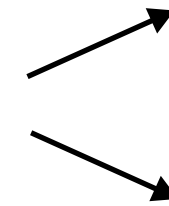


pulse duration =  $0.648 \Delta T$  (sech<sup>2</sup> shape)

# Interferometric autocorrelation



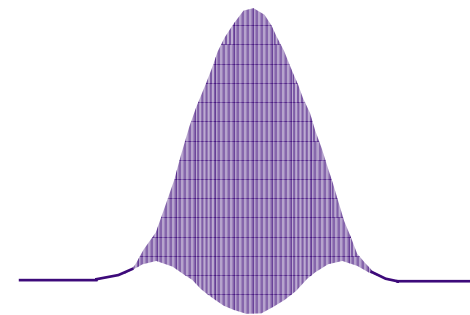
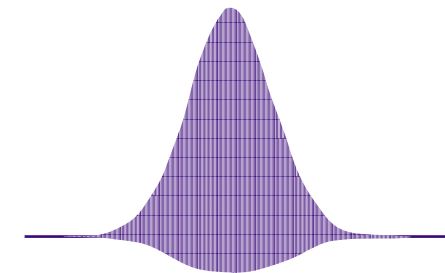
transform-limited



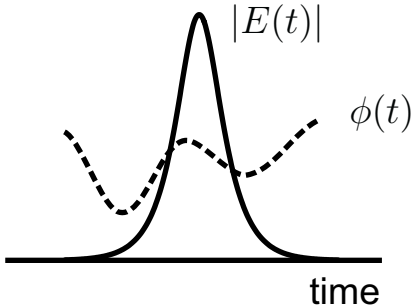
chirped



Interferometric autocorrelation trace



# Characterization of femtosecond pulses

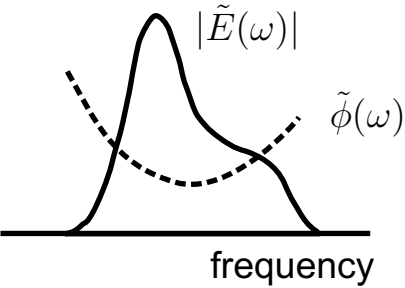


$$E(t) = |E(t)|e^{i\phi(t)}$$



Hard to analyze in femtosecond time scale.

Fourier transformation

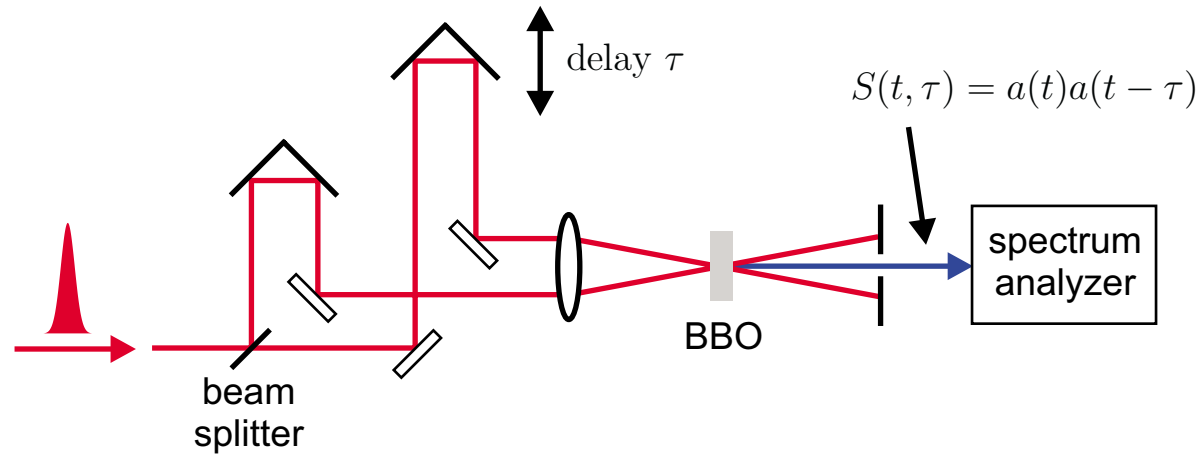


$$\tilde{E}(\omega) = \frac{|\tilde{E}(\omega)|e^{i\tilde{\phi}(\omega)}}{?}$$

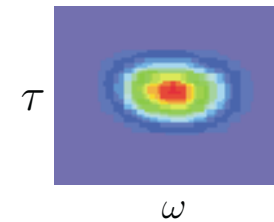


Easy to measure by spectrum analyzer.

# Frequency-resolved optical gating



spectrogram  $|\tilde{S}(\omega, \tau)|_M^2$



measured data

initial guess

$$\tilde{S}(\omega, \tau) = |\tilde{S}(\omega, \tau)|_M e^{i\text{Arg}(\tilde{S})}$$

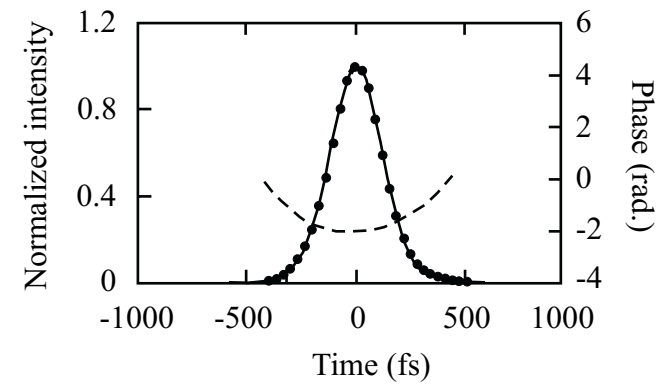
$$\xrightarrow{\text{F.T.}^{-1}} S(t, \tau)$$

replace  
 $|\tilde{S}(\omega, \tau)|$  by  $|\tilde{S}(\omega, \tau)|_M$

$$a(t) = \int S(t, \tau) d\tau$$

$$\tilde{S}(\omega, \tau) = |\tilde{S}(\omega, \tau)| e^{i\text{Arg}(\tilde{S})} \xleftarrow{\text{F.T.}}$$

$$S(t, \tau) = a(t)a(t - \tau)$$



D. J. Kane and R. Trebino, *Opt. Lett.* **18**, 823 (1993).

All electronic phase retrieval measurement using  
nonlinear detectors 📄

1. two-photon interferometry

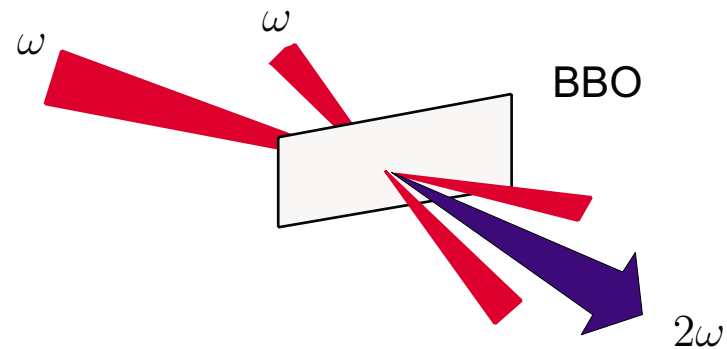
2. waveform reconstruction algorithm



## Limitations of optical nonlinear wave mixing



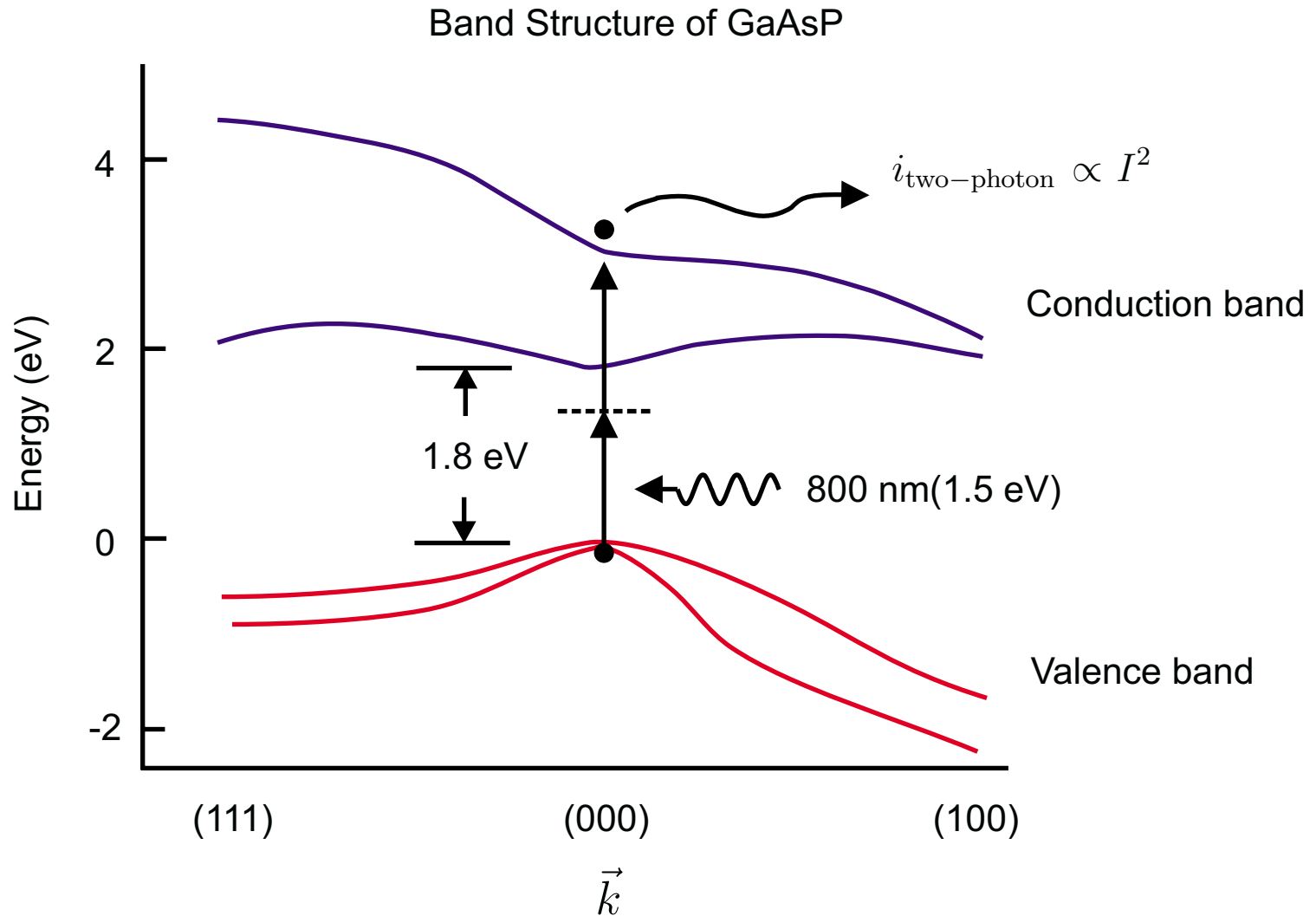
nonlinear wave mixing



Limitations:

1. Dispersion: causes pulse broadening and walk-off.
2. Phase-matching bandwidth: requires thin crystal.
3. Inefficient process: limits sensitivity.

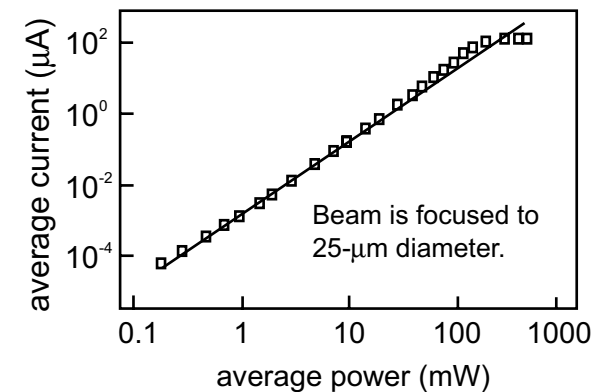
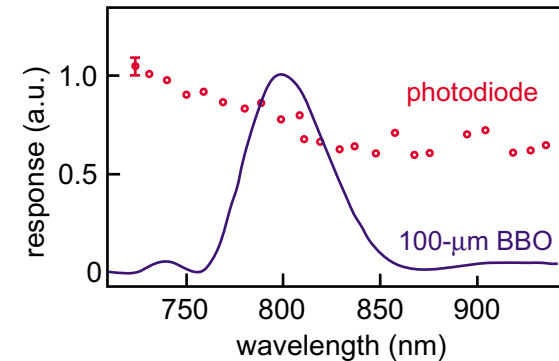
# Two-photon induced photocurrent in semiconductor



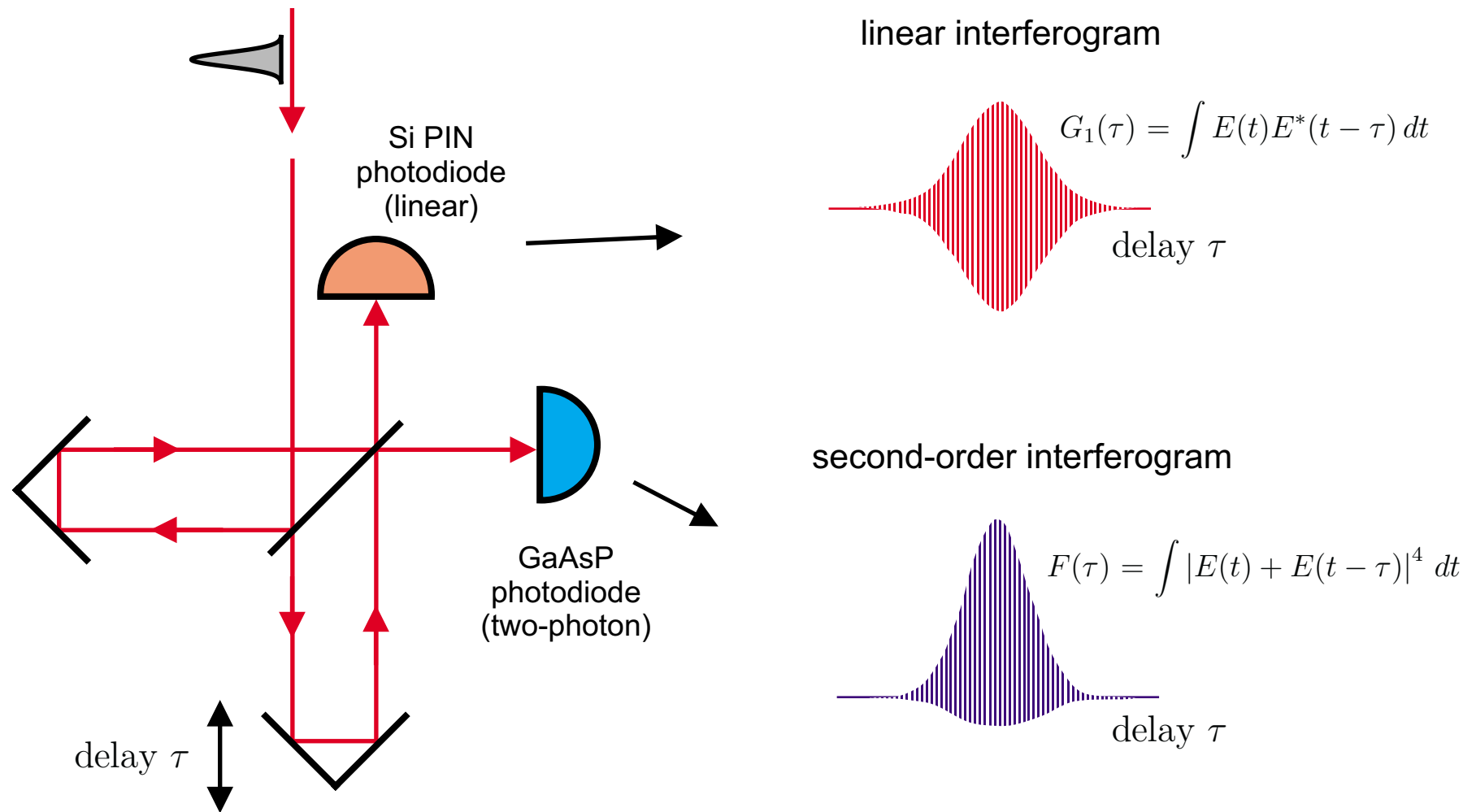
## Advantages of using two-photon photodiodes for nonlinear measurements



1. Bandwidth: nearly uniform up to 200 nm. →
2. Sensitivity: up to  $1.42 \text{ nA/mW}^2$ , comparable to 100- $\mu\text{m}$  BBO + photomultiplier. →
3. Simplicity: easy alignment and no phase-matching requirements.
4. Low cost: less than US\$100.



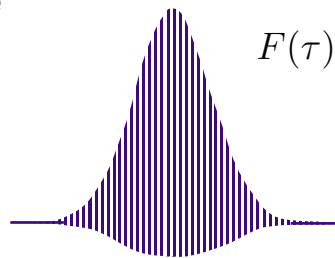
# Two-photon photodiode-based phase-retrieval autocorrelator



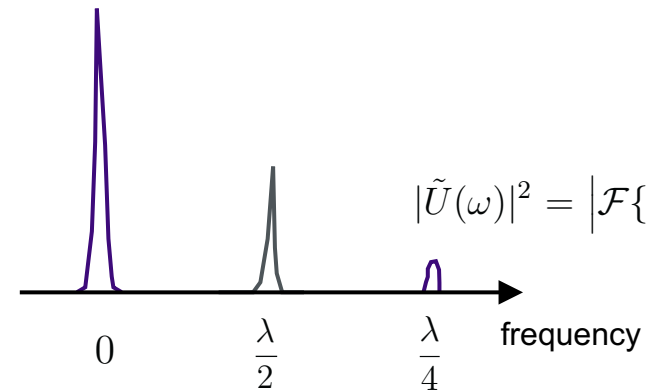
# Spectral components of the autocorrelation function



from two-photon  
photodiode

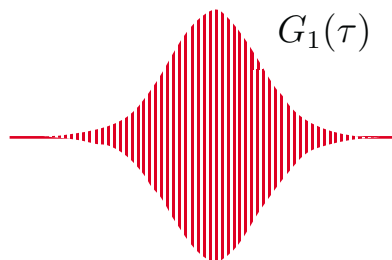


$$|\tilde{I}(\omega)|^2 = |\mathcal{F}\{I(t)\}|^2$$

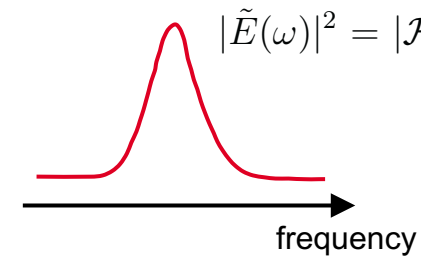


Fourier  
transform,  $\mathcal{F}$

from linear  
photodiode



$$|\tilde{E}(\omega)|^2 = |\mathcal{F}\{E(t)\}|^2$$



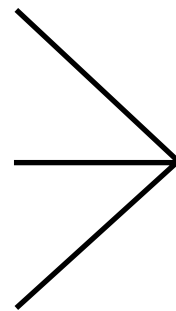
## Unique phase retrieval from measured spectrums



$$|\tilde{E}(\omega)|^2 = |\mathcal{F}\{E(t)\}|^2$$

$$|\tilde{I}(\omega)|^2 = |\mathcal{F}\{I(t)\}|^2$$

$$|\tilde{U}(\omega)|^2 = |\mathcal{F}\{E^2(t)\}|^2$$



uniquely determine the waveform

Ref. : K. Naganuma, K. Mogi, and H. Yamada, *IEEE J. Quantum Electron.* **25**, 1225 (1989).

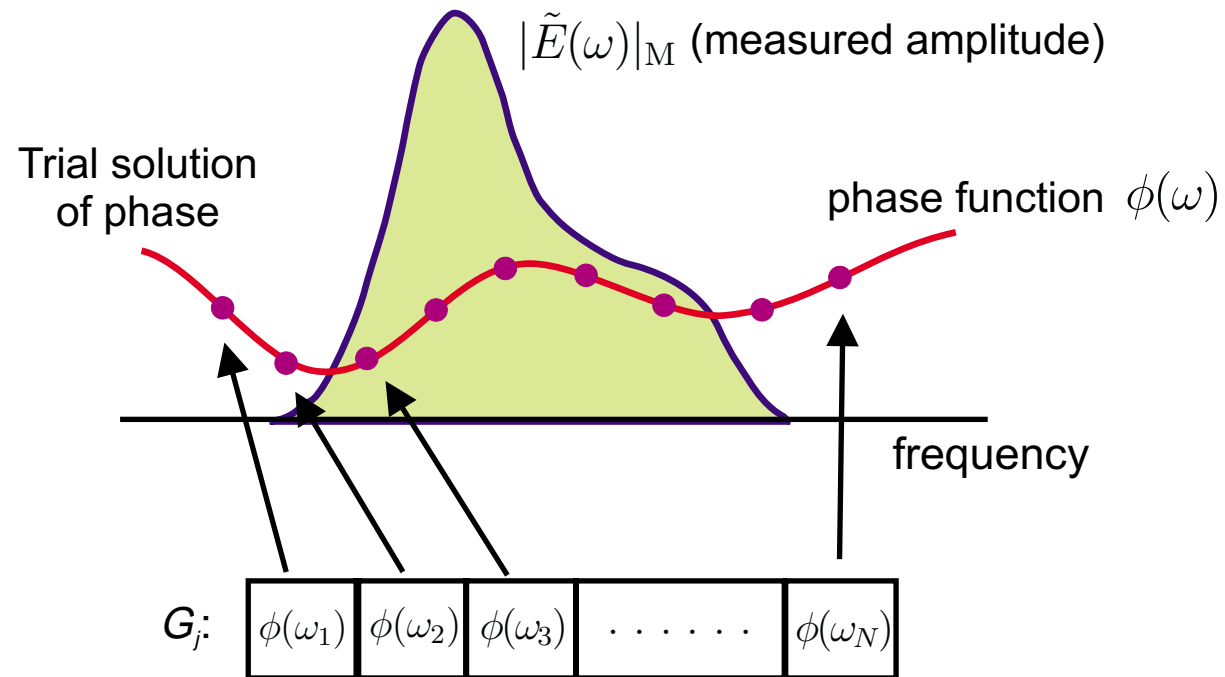
# Genetic Algorithm

1. crossover

2. mutation

3. selection

# Representing the phase function by genes

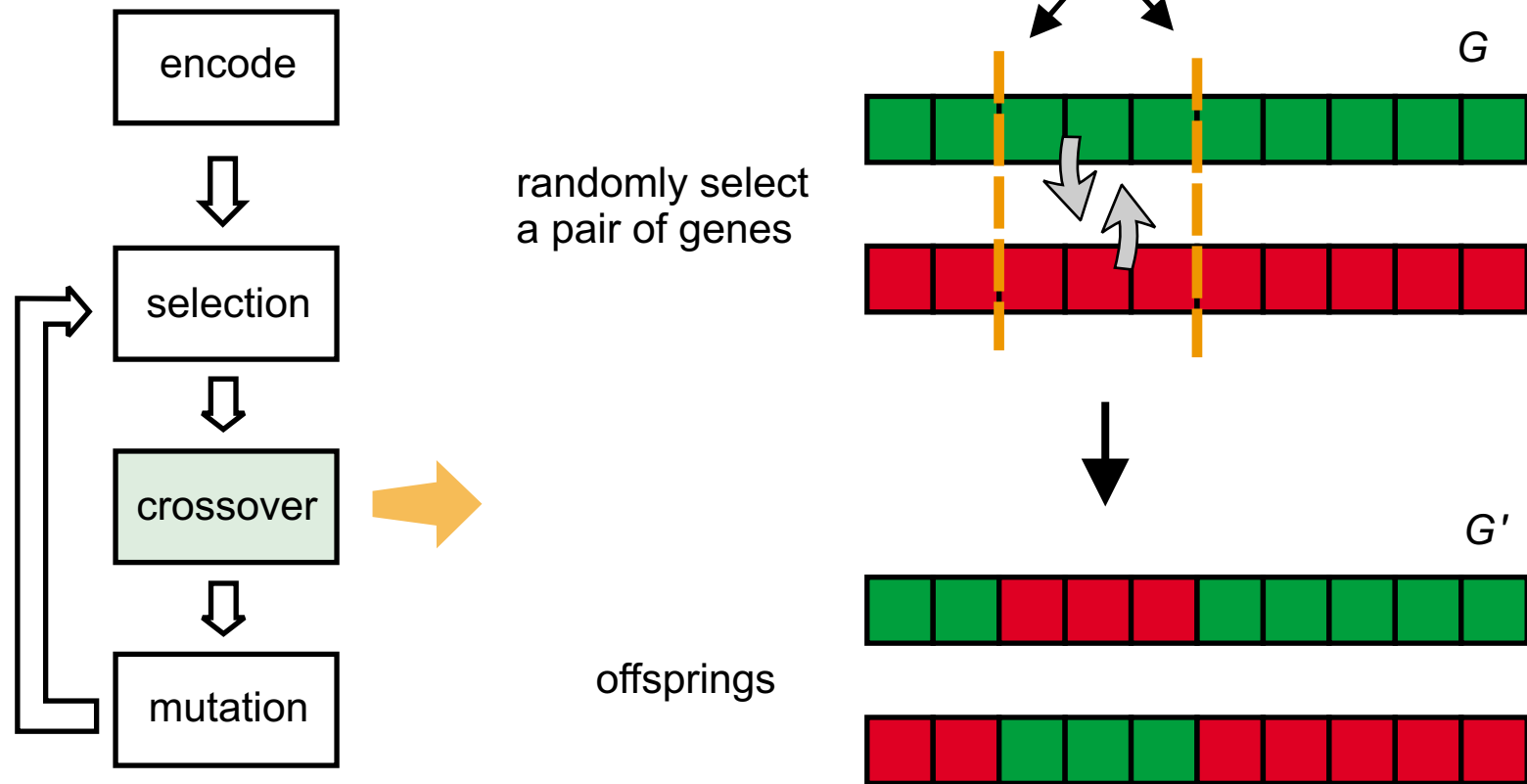


$$\tilde{E}(\omega) = |\tilde{E}(\omega)|_M \cdot \exp[i\phi(\omega)]$$

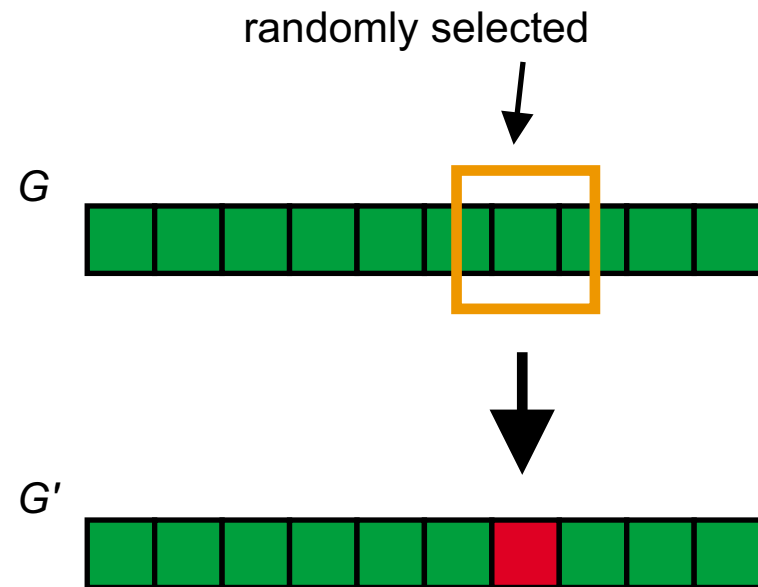
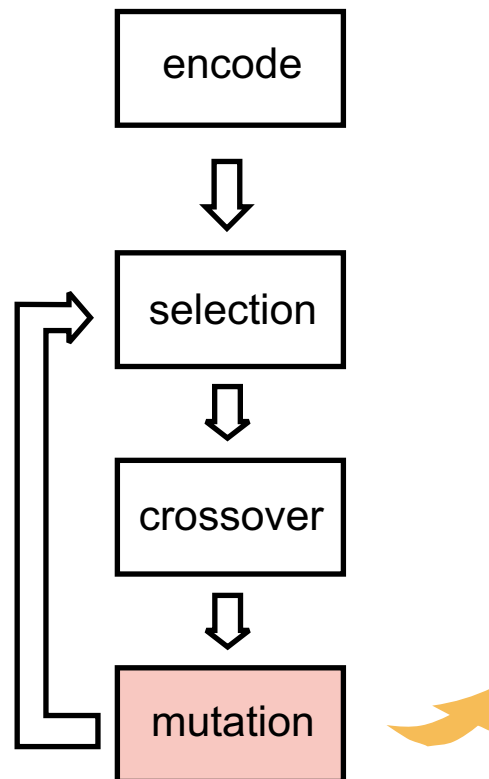
$$E(t) = \mathcal{F}^{-1}\{\tilde{E}(\omega)\}$$



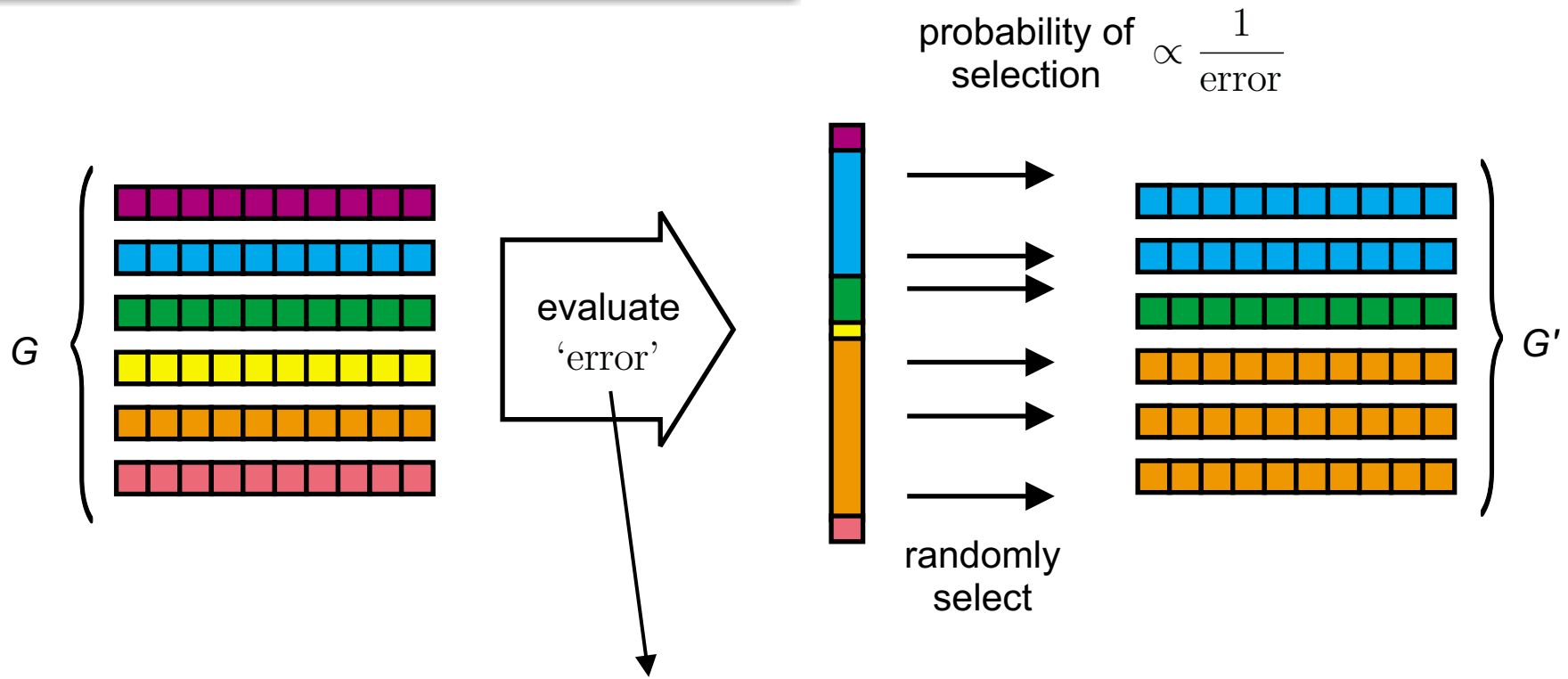
# Basic operation: crossover



## Basic operation: mutation



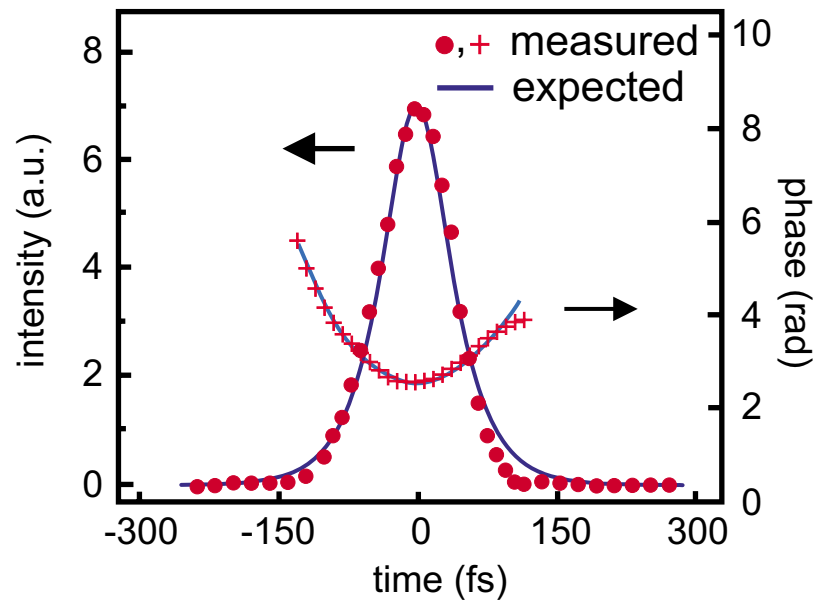
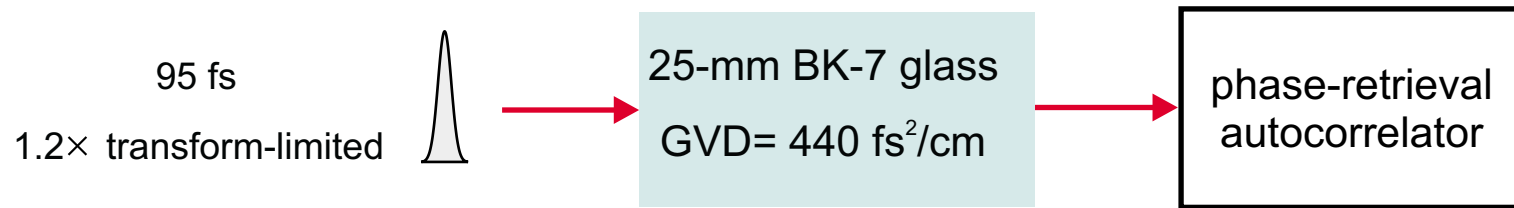
# Basic operation: selection



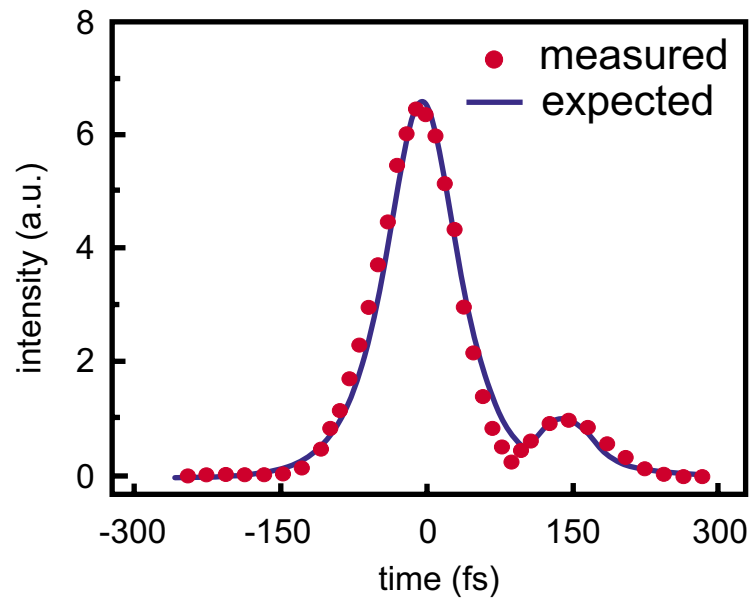
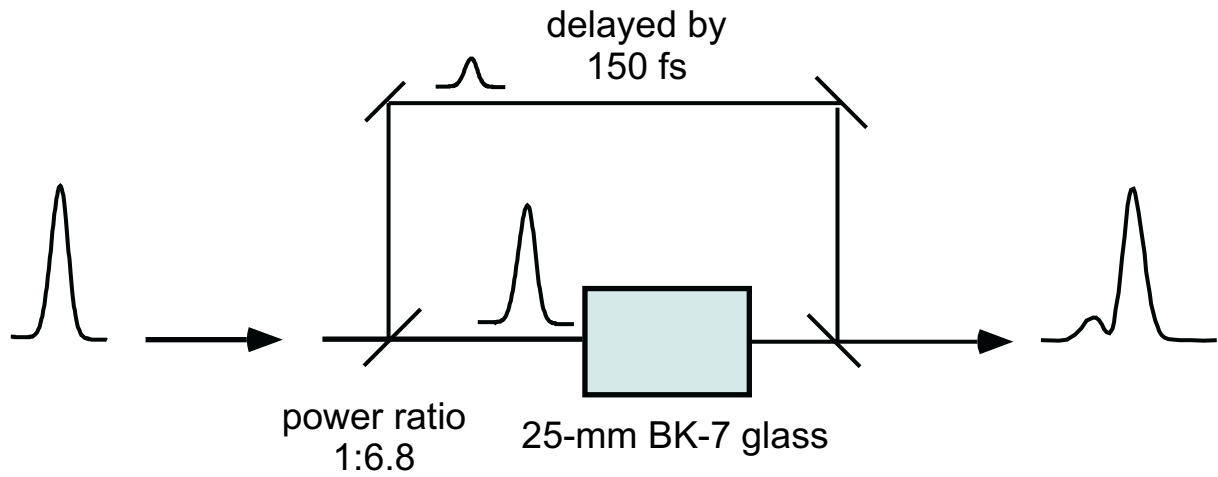
$$E(t) \begin{cases} \rightarrow I(t) = E^*(t)E(t) \xrightarrow{\mathcal{F}} \tilde{I}(\omega) \\ \rightarrow U(t) = E(t)E(t) \xrightarrow{\mathcal{F}} \tilde{U}(\omega) \end{cases}$$

$$\text{error} = \sqrt{Z_I^2 + Z_U^2}, \text{ where } Z_I^2 = \frac{\sum(|\tilde{I}(\omega)|^2 - |\tilde{I}(\omega)|_M^2)^2}{\sum |\tilde{I}(\omega)_M|^4} \text{ and } Z_U^2 = \frac{\sum(|\tilde{U}(\omega)|^2 - |\tilde{U}(\omega)|_M^2)^2}{\sum |\tilde{U}(\omega)_M|^4} .$$

# Demonstration 1: chirped pulses

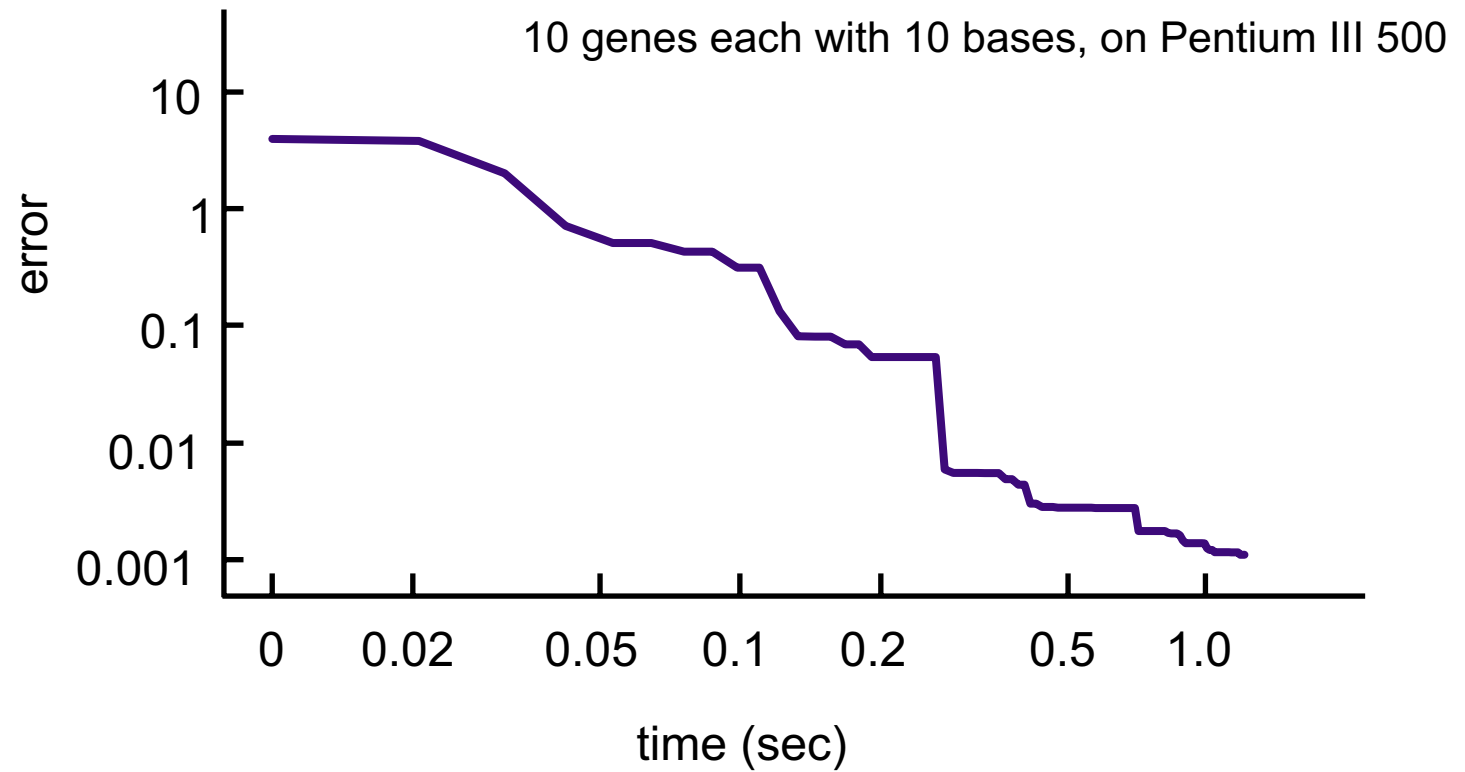


## Demonstration 2: double-peak pulse

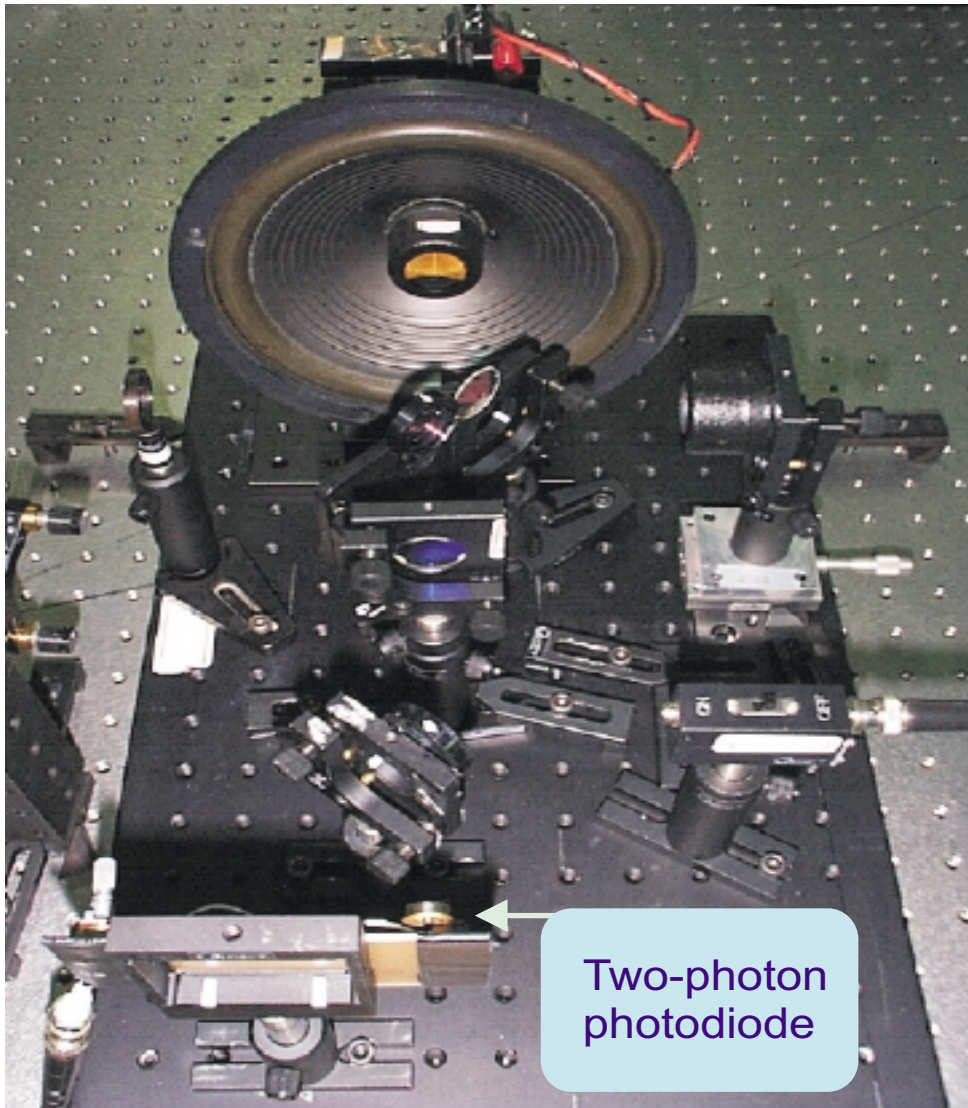


intensity profile of a double-peak pulse

## Converging rate in phase retrieval with genetic algorithm



## Semiconductor Phase-Retrieval Autocorrelator



45cm\*30cm\*30cm

## Potential applications

1. ultrafast spectroscopy

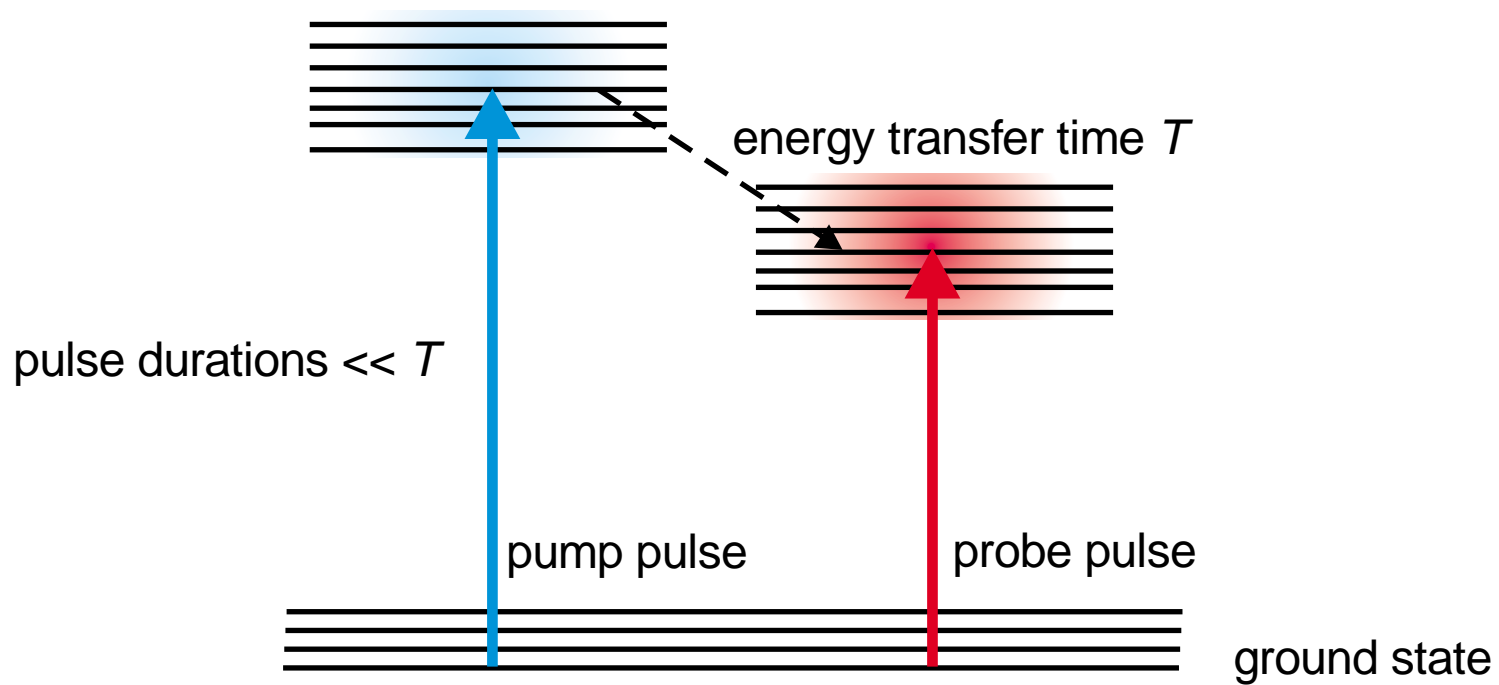
2. accelerator beam diagnosis



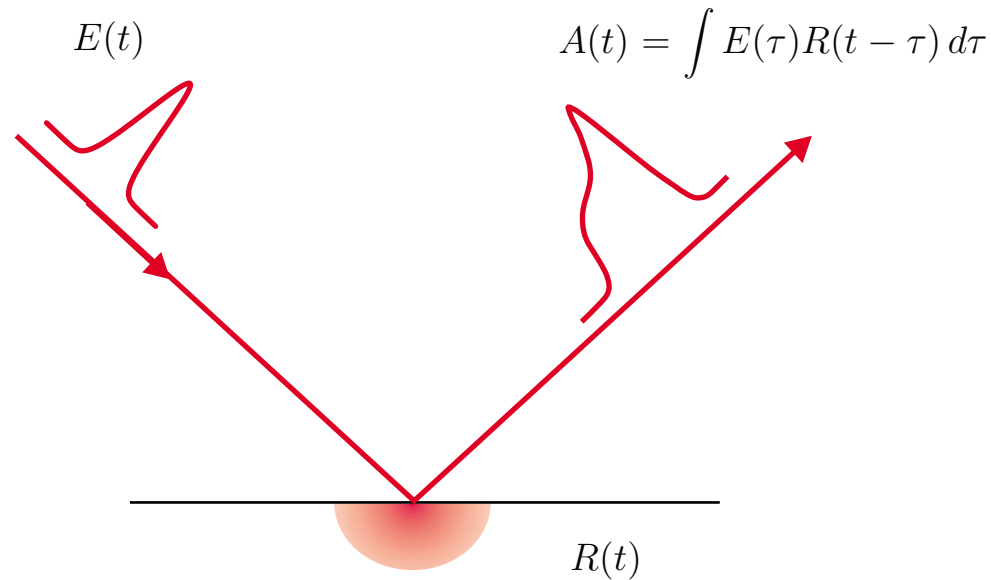
# Ultrafast pump-probe experiments



## transient absorption experiment

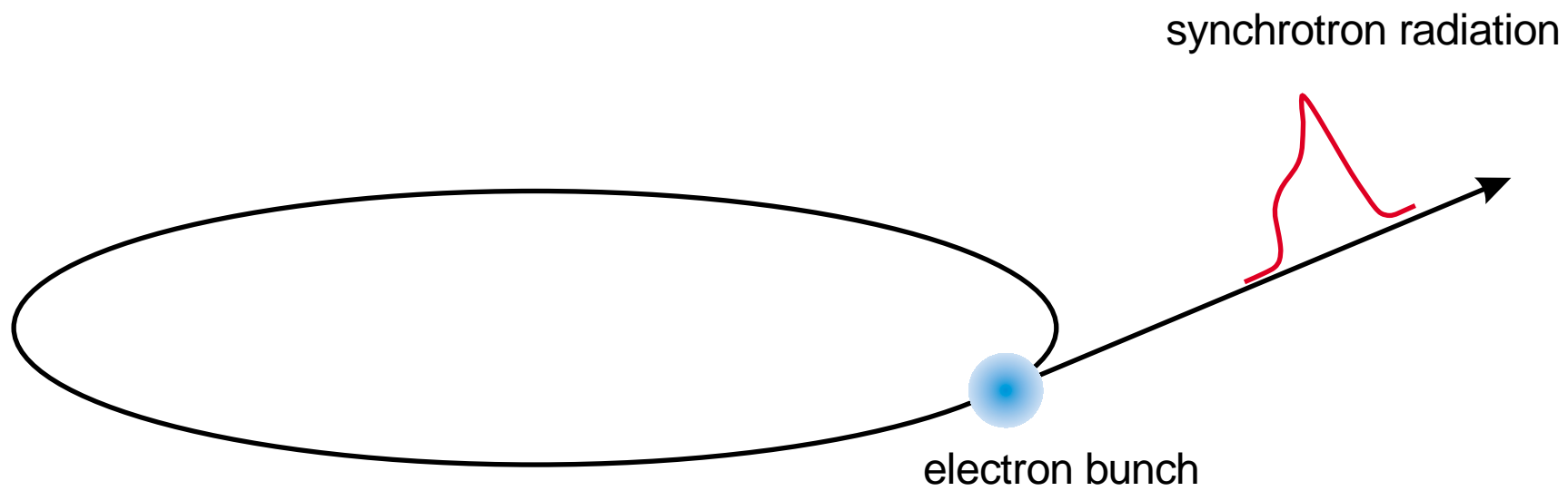


## Avoiding excessive excitation bandwidth






It is possible to measure the material response with pulse duration of the same time scale, thereby **avoiding excessive excitation bandwidth**.

# Beam diagnosis



## Summary



-  Slow nonlinear detectors can measure fast signals, without nonlinear wave mixing.
-  From nonlinear interference, complete waveform can be reconstructed by the genetic algorithm.
-  The technique is demonstrated with ultrafast optics, and can be used in other fields.